

So going back to basics, the mode of bending correction used on the x26-c mono is to apply a torque to one end of a plate with the other end fixed. this make a linear increase in angle from the fixed point to the end of the second crystal. the angle is described by.  $\Delta\phi(x) = (\frac{ML}{EI})x$  where  $M = appliedTorque$  and  $L = lengthofsecondcrystal$ ,  $I$  is thin plate inertia as  $I = \frac{1}{12}m(L^2 + w^2)$  and  $E$  is the Elastic modulus of Si. This gives a parabolic displacement by  $\Delta y(x) = \frac{Mx^2}{2EI}$

so that with a non zero corrective bend applied the reflecting condition of the second crystal is energy dependant. with *crystal1* set at  $\theta$  with the  $Si[111]$  reflection selecting the energy by  $\lambda = 2\frac{d}{[111]} \sin \theta$  and the second crystal selecting energy by  $\lambda = 2\frac{d}{[111]} \sin(\theta + \Delta\phi)$  since  $\Delta\phi$  is set by  $H/\tan \theta = x$  so that  $\Delta\phi = (\frac{ML}{EI})H\frac{1}{\tan \theta}$  where  $H$  is the inter-crystal spacing.

$$\lambda = 2\frac{d}{[111]} \sin(\theta + (\frac{ML}{EI})H\frac{1}{\tan \theta})$$

$$\lambda_2 = 2\frac{(d + \frac{-yLd}{EI}M)}{[111]} \sin(\theta + \frac{[\frac{(\frac{H}{\tan \theta})^2 M}{2EI}] + H}{\tan \theta})$$

Using a calculated Corrective motion the scanning time to find the maximum intensity for a given energy may be greatly reduced. Because of indirect coupling of the motor motion to the bending moment arm, and motor slippage under load it is not possible to move directly to the correct position by calculation. Through use of a simple strain gauge bonded to the non reflecting side of the second crystal closed feed back for corrective bending could be accomplished.

Thermal expansion will also play a role in observed intensity, as the first crystal heats it will begin to expand along a line at its middle where the beam falls on it. the expansion will cause a local increase in the first crystal's  $\theta$ . and a translation further towards the source point. both effects though relatively small will translate the reflection towards  $x = 0$  on the second crystal. which will reduce the  $\Delta\phi$  if this reduction is faster than the local increase in  $\theta$  on the first crystal the reflecting condition will be improved. if the local increase in  $\theta$  on the first crystal is greater the error will be magnified.

There is also an error due to change in  $d$  spacing of the primary crystal it is governed by the Debye temperature factor.

$$2B = \frac{12h^2T}{m_a k \Theta_D^2} Q(\frac{\Theta_D}{T})$$

Look up  
the relation  
of 2B to  
D-spacing,  
it should be  
in azsaroff

where the debye tempature is set by  $\nu_m = (\frac{3N}{4\pi V})^{\frac{1}{3}}v_s$  and  $\Theta_D = \frac{h\nu_m}{k}$  where  $N/V$  is the number density of atoms and  $v_s$  is the effective speed of sound in the solid

Thermal expansion is defined by  $\alpha = \frac{\gamma C_V}{K_T V} = \frac{\gamma C_P}{K_S V}$  where  $\gamma$  is the heat capacity ratio,  $C_V$  is the heat capacity at constant volume,  $K_T$  is the isothermal bulk modulus,  $C_P$  is the heat capacity at constant pressure, and  $K_S$  is the adiabatic bulk modulus and  $V$  is the volume also  $\alpha = 3\beta$  where  $\beta$  is the liniar expansion.

there is an other mode of change in the D-spacing caused by the bending of the second crtystal. related to the radius of curverture  $\rho$  which deforms the lattice to a larger D-space.

So for the second crystal